



Aalto University
School of Electrical
Engineering

Scattering from a sphere of nonlinear material

Luis R.J. Costa

*Department of Radio Science and Engineering
Aalto University School of Electrical Engineering*

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Why study scattering from a sphere of nonlinear material?

Model of the problem to be simulated

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Why study scattering from a sphere of nonlinear material?

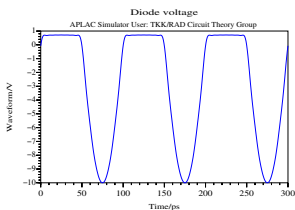
- Nonlinearity generates harmonics

Why study scattering from a sphere of nonlinear material?

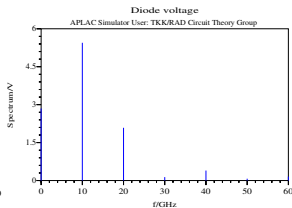
- Nonlinearity generates harmonics
 - How much energy is shifted from the fundamental frequency to the harmonics?

Why study scattering from a sphere of nonlinear material?

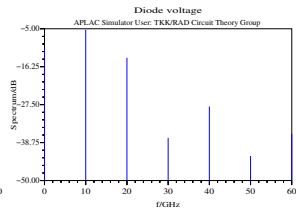
- Nonlinearity generates harmonics
 - How much energy is shifted from the fundamental frequency to the harmonics?
 - Type of nonlinearity used: that of a pn-junction (exp. function)
 - A real material; it has the most extreme nonlinearity possible



Waveform



Spectrum (V)



Spectrum (dB)

Why study scattering from a sphere of nonlinear material?

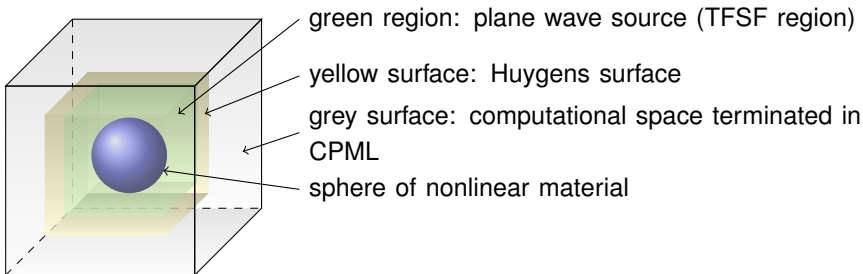
- Why sphere?
 - it is a canonical benchmark whose radar cross-section (RCS) is well known (Mie scattering)

Why study scattering from a sphere of nonlinear material?

- Why sphere?
 - it is a canonical benchmark whose radar cross-section (RCS) is well known (Mie scattering)
- Can nonlinear materials be exploited in stealth technology?

Model of the problem to be simulated

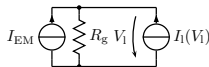
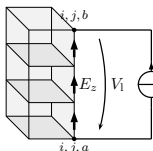
Problem is simulated using the FDTD method



The Huygens surface is required to obtain the far field

Modelling the nonlinear material

Modelling a lumped element in FDTD



$$\oint_C \mathbf{H}^{n+\frac{1}{2}} \cdot d\mathbf{l} = \int_S \left(\epsilon \frac{\partial \mathbf{E}^{n+\frac{1}{2}}}{\partial t} + \sigma \mathbf{E}^{n+\frac{1}{2}} + \mathbf{J}_1^{n+\frac{1}{2}} \right) \cdot d\mathbf{s}$$

Ampere's law

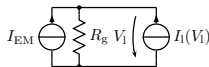
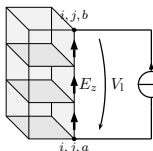
Update equation:

$$E_z|_{i,j,k+\frac{1}{2}}^{n+1} = \left(\frac{1 - \frac{\sigma \Delta t}{2\epsilon}}{1 + \frac{\sigma \Delta t}{2\epsilon}} \right) E_z|_{i,j,k+\frac{1}{2}}^n + \left(\frac{\Delta t}{\epsilon} \right) \times \left\{ \frac{\mathcal{I}_z|_{i,j,k}^n}{\Delta_x \Delta_y} - \frac{I_1(V_1^{n+\frac{1}{2}})}{\Delta_x \Delta_y} \right\}$$

$$\mathcal{I}_z|_{i,j,k}^n = (H_y|_{i+\frac{1}{2},j,k+\frac{1}{2}}^{n+\frac{1}{2}} - H_y|_{i-\frac{1}{2},j,k+\frac{1}{2}}^{n+\frac{1}{2}}) \Delta_y - (H_x|_{i,j+\frac{1}{2},k+\frac{1}{2}}^{n+\frac{1}{2}} - H_x|_{i,j-\frac{1}{2},k+\frac{1}{2}}^{n+\frac{1}{2}}) \Delta_x$$

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Modelling a lumped element in FDTD



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Lumped elements can be directed only in the direction of the primary axes

Modelling the nonlinear material

Modelling a diode spanning one or many Yee cells

Diode current-voltage equation: $I_d = I_s \left(e^{\alpha V_d} - 1 \right)$, $\alpha = q/(kT)$

Modelling the nonlinear material

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■ Diode voltage V_d :

- spanning one cell from i, j, a to $i, j, a + 1$: $V_d^{n+\frac{1}{2}} = -E_z|_{i,j,a}^{n+\frac{1}{2}} \Delta_z$
- spanning cells from i, j, a to i, j, b : $V_d^{n+\frac{1}{2}} = -\Delta_z \sum_{k=a}^{b-1} E_z|_{i,j,k}^{n+\frac{1}{2}}$

Modelling the nonlinear material

Modelling a diode spanning one or many Yee cells

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- Iterate $V_d^{n+\frac{1}{2}}$ using, e.g., Newton-Raphson:

$$V_d^{n+\frac{1}{2},l+1} = V_d^{n+\frac{1}{2},l} - \frac{V_d^{n+\frac{1}{2},l}/R_g + I_s \left(e^{\frac{q}{kT} V_d^{n+\frac{1}{2},l}} - 1 \right) - I_{EM}}{1/R_g + \frac{q}{kT} I_s e^{\frac{q}{kT} V_d^{n+\frac{1}{2},l}}}, \quad V_d^{n+\frac{1}{2},0} = V_d^{n-\frac{1}{2}}$$

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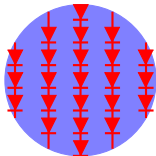
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- Compute I_d and plug in $I_l = -I_d$ via the update equations in the column of cells in which it exists

Sphere of nonlinear material

- Sphere is filled with z-directed diodes
- A diode current passes through every Yee cell within the sphere



Encountered problems and work-arounds

- The presence of many diodes in close proximity makes the simulation susceptible to instability
 - the material is strongly nonlinear
- The time step may have to be reduced from the largest possible step obtained via the Courant-Friedrichs-Lewy condition
- The plane-wave amplitude cannot be arbitrarily large
 - the strong nonlinearity in the current-voltage dependence causes instability
 - the larger the field strength, the larger the voltage implying larger changes in current even for relatively small changes in the large voltage

Simulations

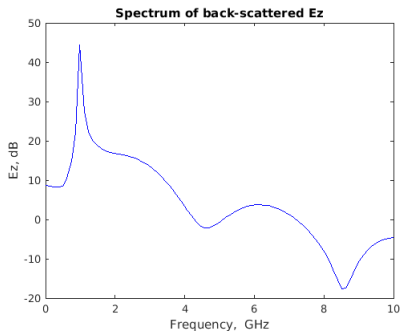
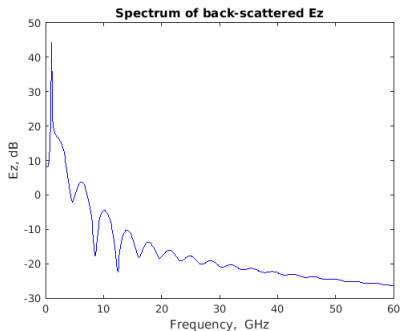
- Different spheres are simulated for
 - wavelength much larger than the sphere
 - wavelength of similar size as the sphere

Simulation parameters

- Grid size: $120 \times 120 \times 120$
- Discretisation: $\Delta_x = \Delta_y = \Delta_z = 0.5$ mm,
time step damping factor: $0.5 \Rightarrow \Delta_t = 0.963$ ps
- Plane wave propagates along x-axis with amplitude 5000 V/m
- Diode saturation current: 1 pA (typical value for a silicon diode)
- Sphere diameter: 32 mm (64 cells at equatorial plane)
- Scattered-field observation point: 15, 60, 60 (between CPML and the Huygens surface)

Simulations: first a PEC sphere

Case 1a: $f = 1$ GHz, $\lambda \approx 300$ mm, time steps $n = 16\,384$

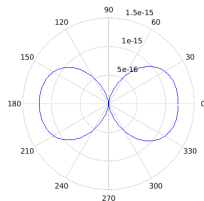
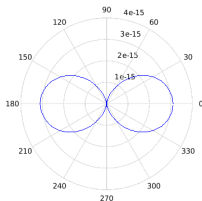
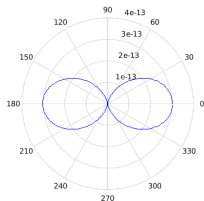


- The plane-wave frequency is obvious.
- The model errors are also obvious
 - the error can be reduced but not eliminated by increasing n

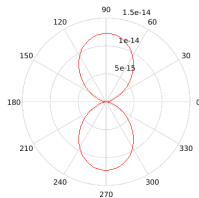
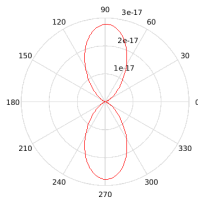
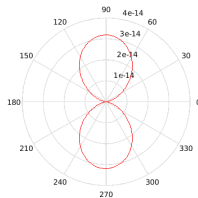
Simulations: first a PEC sphere

Case 1a: $f = 1 \text{ GHz}$, $\lambda \approx 300 \text{ mm}$ (arbitrary scale)

RCS_{θ}



RCS_{ϕ}



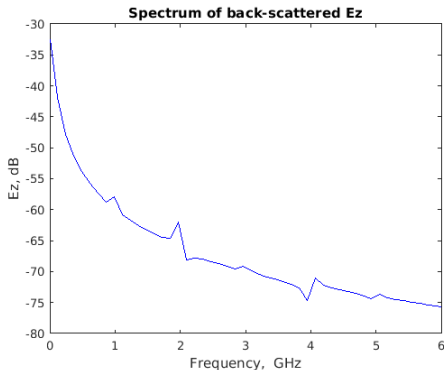
f

$2f$

$3f$

Simulations: diode sphere

Case 1b: $f = 1$ GHz, $\lambda \approx 300$ mm, $n = 16\,384$, $\epsilon_r = 1$

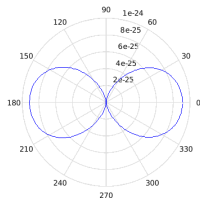
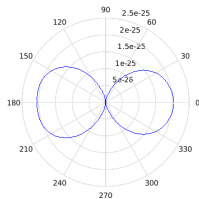
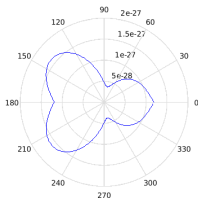


- A very strong DC component (for electromagnetic rectification) is present
- Harmonic peaks are distinguishable

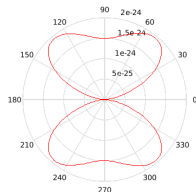
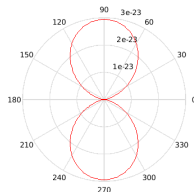
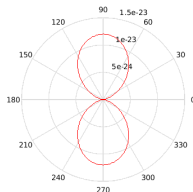
Simulations: diode sphere

Case 1b: $f = 1 \text{ GHz}$, $\lambda \approx 300 \text{ mm}$, $\epsilon_r = 1$ (arbitrary scale)

RCS_θ



RCS_ϕ



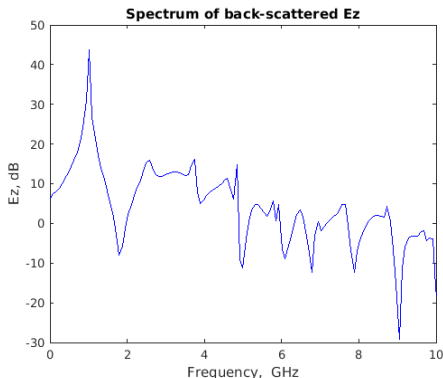
f

$2f$

$3f$

Simulations: diode sphere

Case 1c: $f = 1 \text{ GHz}$, $\lambda \approx 300 \text{ mm}$, $n = 16384$, $\epsilon_r = 11.68$

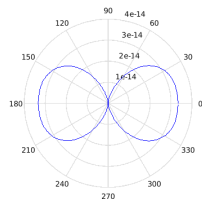
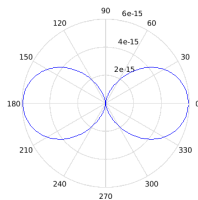
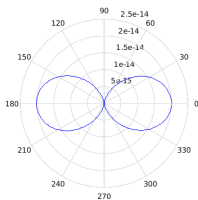


- The fundamental frequency is clearly visible

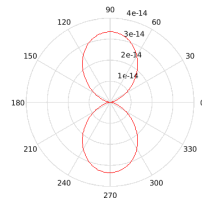
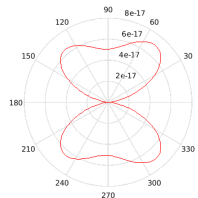
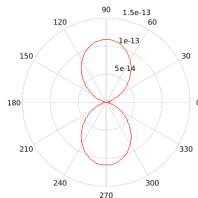
Simulations: diode sphere

Case 1c: $f = 1$ GHz, $\lambda \approx 300$ mm, $\epsilon_r = 11.68$ (arbitrary scale)

RCS_θ



RCS_ϕ



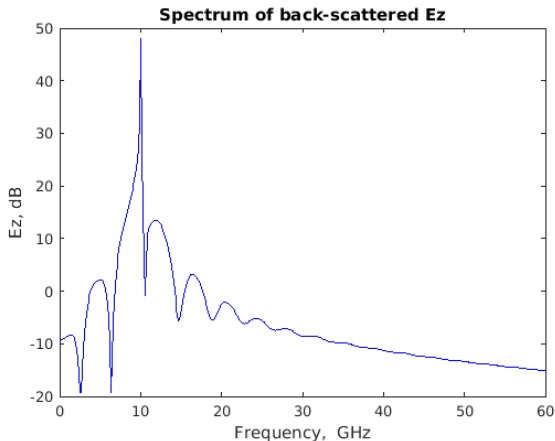
f

$2f$

$3f$

Simulations: again a PEC sphere

Case 2a: $f = 10$ GHz, $\lambda \approx 30$ mm, $n = 16384$

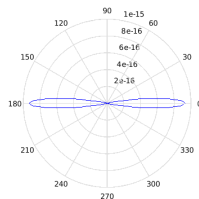
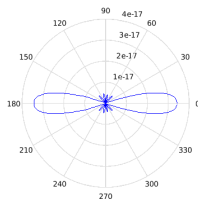
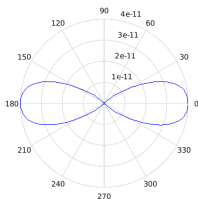


The plane-wave frequency and model error are obvious here too.

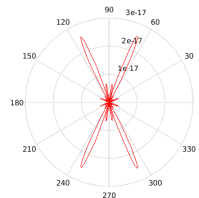
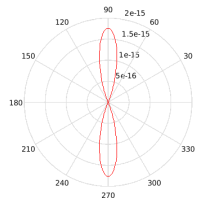
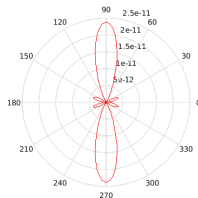
Simulations: PEC sphere

Case 2a: $f = 10$ GHz, $\lambda \approx 30$ mm (arbitrary scale)

RCS_{θ}



RCS_{ϕ}



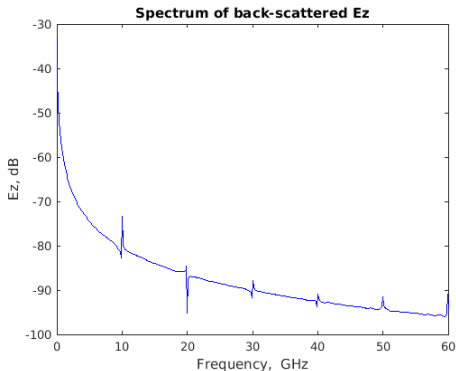
f

$2f$

$3f$

Simulations: diode sphere

Case 2b: $f = 10 \text{ GHz}$, $\lambda \approx 30 \text{ mm}$, $n = 16\,834$, $\epsilon_r = 1$

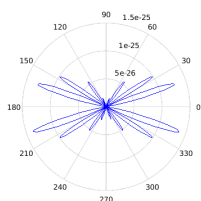
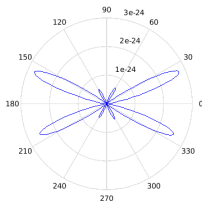
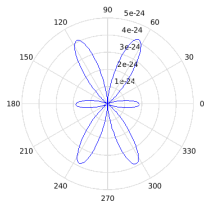


- A very strong DC component is present
- The fundamental frequency and harmonics are visible

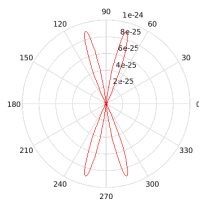
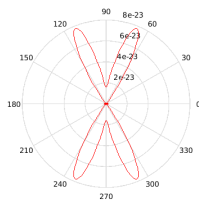
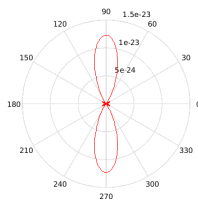
Simulations: diode sphere

Case 2b: $f = 10$ GHz, $\lambda \approx 30$ mm, $\epsilon_r = 1$ (arbitrary scale)

RCS_{θ}



RCS_{ϕ}



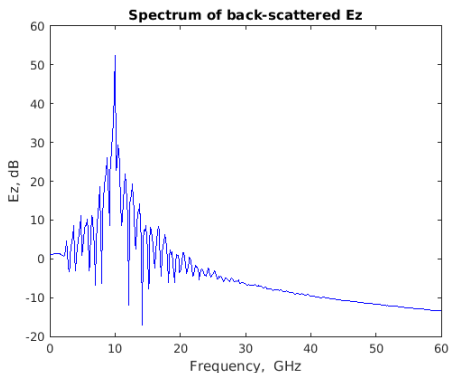
f

$2f$

$3f$

Simulations: diode sphere

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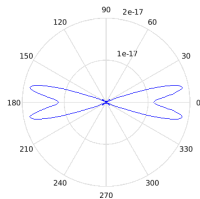
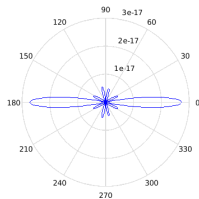
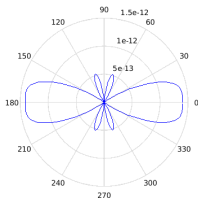


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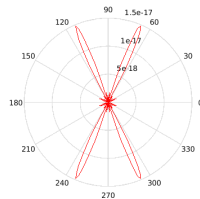
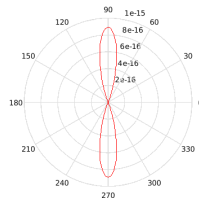
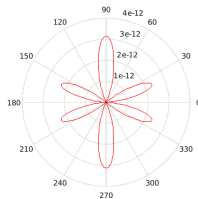
Simulations: diode sphere

Case 2c: $f = 10$ GHz, $\lambda \approx 30$ mm, $\epsilon_r = 11.68$ (arbitrary scale)

RCS_θ



RCS_ϕ



f

$2f$

$3f$

What do the data tell us?

$$\epsilon_r = 1$$

- Harmonics are generated as expected, but they are weak
- Charge accumulates at the nodes of the diodes since they have no discharge path available
 - the scattered field stays positive in the back-scattered direction
 - a prominent DC component emerges in the spectra as a result

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- Effects of nonlinearity are no longer apparent
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$$\epsilon_r = 11.68$$

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Other remarks

- Several cycles of the plane wave are required for the nonlinearities to kick in
- A relatively wide modulated pulse is necessary for the nonlinearities to have an effect

What next?

- Is it possible to arrange the diodes so as to enhance the effects of nonlinearities?
- Can the incident polarisation current be manipulated through a suitable arrangement of the diodes?
- Study how different pulses interact with spheres of such nonlinear materials

Conclusions

- A model for a nonlinear material is successfully implemented and used to compute scattering from a sphere made of a nonlinear material
- Simulation results show interesting behaviour in the sphere modelled
 - harmonics are created, but not always
- Energy is shifted primarily to DC, but to other frequencies as well
- Application to stealth technology requires further study